Entanglement entropies in conformal systems with boundaries

Luca Taddia

Scuola Normale Superiore, Pisa & CNR-INO, Sesto Fiorentino

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Analitical results Boundary CFT Entanglement entropies and BCFT

Numerical tests Ising chain XX chain

Work in collaboration with:





German Sierra (IFT UAM-CSIC)

Jose C. Xavier (Instituto de Física, Universidade Federal de Uberlândia) Francisco C. Alcaraz (Instituto de Física de São Carlos,

Universidade de São Paulo)

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Why study entanglement in many body physics?

 Understanding correlations is hard: one needs non-local quantities in order to fully get what is happening in the system Entanglement entropies in conformal systems with boundaries

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- Strictly related topic: efficiency of numerical methods based on matrix-product states (often numerics is the most reliable way for extracting theoretical information from quantum Hamiltonians)

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- Most studied entanglement witnesses: Rényi entanglement entropies:

$$S_n(A) = \frac{1}{1-n} \mathrm{Tr}_A \rho_A^n$$

 $\rho_{A}:$ reduced density matrix of subsystem A

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► As $n \to 1$: von Neumann entanglement entropy $S = -\text{Tr}_A \left[\rho_A \ln \rho_A \right]$ Entanglement entropies in conformal systems with boundaries

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Rényi entropies and CFT

Seminal result [Holzhey et al. '94]: logarithmic violation of the area law: for a finite interval in an infinite 1D chain,

$$S = \frac{c}{3} \ln l$$

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Rényi entropies and CFT

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$$S = \frac{c}{3} \ln b$$

 Generalizations at finite size, PBC/FBC and generic n [Calabrese & Cardy '04]:

$$S_n = rac{c}{6\eta} \left(1 + rac{1}{n}
ight) \ln \left[rac{\eta L}{\pi} \sin rac{\pi I}{L}
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with $\eta=$ 1, 2 for PBC/OBC

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with $\eta=$ 1, 2 for PBC/OBC

These results were checked numerically in a plenty of models, and constitute now the easiest way of extracting c from simulations Entanglement entropies in conformal systems with boundaries

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$$\bullet \ \rho_A^n = \underbrace{\rho_A \times \rho_A \times \cdots \times \rho_A}_n$$

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$$\rho_A^n = \underbrace{\rho_A \times \rho_A \times \cdots \times \rho_A}_n$$

Replica trick: computing ρⁿ_A is equivalent to computing the reduced density matrix on a particular Riemann manifold R_n

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Taking the trace (i.e., computing the n-th Rényi entropy) amounts to sew together the last and the first sheet Entanglement entropies in conformal systems with boundaries

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 If the system is conformal invariant, the space time can be mapped to simpler geometries by means of appropriate conformal transformations (operators also transform according to them) Entanglement entropies in conformal systems with boundaries

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For the ground state, Trρⁿ_A is seen to behave, under conformal transformations, as a two-point function of primary-like operators (twist fields: see, e.g., [Castro-Alvaredo & Doyon '09]) of conformal dimensions depending on n

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What happens when consider an excited state? [Alcaraz et al. '11, Berganza et al. '12] Entanglement entropies in conformal systems with boundaries

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What happens when consider an excited state? [Alcaraz et al. '11, Berganza et al. '12]

CFT picture: in radial quantization:

$$\left|h, \bar{h}\right\rangle = \lim_{z, \bar{z} o 0} \Upsilon^{\dagger}\left(z, \bar{z}\right) \left|0
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 Υ : primary field of conformal dimensions h, \bar{h}

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► Replica manifold for a closed finite system: Y has to be placed at t = ±∞ on each sheet



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• The excess entropy wrt the GS, $F_{\Upsilon}^{(n)}(x) \equiv \operatorname{Tr}\rho_{A,\Upsilon}^{n}/\operatorname{Tr}\rho_{A,\mathbb{I}}^{n}, \text{ takes the form}$ $F_{\Upsilon}^{(n)}(x) = \lim_{w_{k} \to -i\infty} \frac{\left\langle \prod_{k=0}^{n-1} \Upsilon(w_{k}, \bar{w}_{k}) \Upsilon^{\dagger}(-w_{k}, -\bar{w}_{k}) \right\rangle_{\mathcal{R}_{n}}}{\left\langle \Upsilon(w_{0}, \bar{w}_{0}) \Upsilon^{\dagger}(-w_{0}, -\bar{w}_{0}) \right\rangle_{\mathcal{R}_{1}}^{n}}$

 \blacktriangleright Conformal transformations: the replica manifold is mapped to $\mathbb C$







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► Conformal transformations: the replica manifold is mapped to C



The operators are placed at z[±]_{n,k} = exp [^{iπ}/_n(±x + 2k)], with x ≡ ¹/_I

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► Conformal transformations: the replica manifold is mapped to C



- ► The operators are placed at $z_{n,k}^{\pm} = \exp\left[\frac{i\pi}{n}(\pm x + 2k)\right]$, with $x \equiv \frac{l}{L}$
- The excess entropy is now

$$F_{\Upsilon}^{(n)}(x) = n^{-2n(h+\bar{h})} \frac{\left\langle \prod_{k=0}^{n-1} \Upsilon\left(z_{n,k}^{+}, \bar{z}_{n,k}^{+}\right) \Upsilon^{\dagger}\left(z_{n,k}^{-}, \bar{z}_{n,k}^{-}\right) \right\rangle}{\left\langle \Upsilon\left(z_{1,0}^{+}, \bar{z}_{1,0}^{+}\right) \Upsilon^{\dagger}\left(z_{1,0}^{-}, \bar{z}_{1,0}^{-}\right) \right\rangle^{n}}$$

where the correlators are on \mathbb{C} (\Box) (\Box)

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Question: what happens when one considers open systems?

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Question: what happens when one considers *open* systems? Can one generalize the Calabrese-Cardy formula to excited states or different kind of boundary conditions? Entanglement entropies in conformal systems with boundaries

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General questions: how boundaries affect conformal properties? Do boundary conditions preserving conformal invariance exist? Entanglement entropies in conformal systems with boundaries

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- General questions: how boundaries affect conformal properties? Do boundary conditions preserving conformal invariance exist?
- Answer [Cardy '89]: for rational theories, they exist and they are in one-to-one correspondence to the primary fields of the theory

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- General questions: how boundaries affect conformal properties? Do boundary conditions preserving conformal invariance exist?
- Answer [Cardy '89]: for rational theories, they exist and they are in one-to-one correspondence to the primary fields of the theory
- In general, the partition function of a rational model is given by $Z = \sum_{h,\bar{h}} \mathcal{M}_{h,\bar{h}} \times \chi_h \times \chi_{\bar{h}}$, $\mathcal{M}_{h,\bar{h}}$: multiplicity of the Verma module $h \otimes \bar{h}$

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- Boundaries reduce the operator content of the theory: just one chirality survives, and

$$Z_{\alpha\beta} = \sum_{h} \mathcal{N}^{h}_{\alpha\beta} \chi_{h}$$

 $\mathcal{N}: \text{ fusion coefficients} \rightarrow \text{correspondence between}$ boundary conditions and primary fields Entanglement entropies in conformal systems with boundaries

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Consequence: the GS of a system with boundary conditions α, β is obtained from the GS with free BC (FBC) by applying on it a chiral primary operator Υ (see also H. Saleur's talk)

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- Consequence: the GS of a system with boundary conditions α, β is obtained from the GS with free BC (FBC) by applying on it a chiral primary operator Υ (see also H. Saleur's talk)
- What we have now is a formula very similar to the PBC one for a primary excited state:

$$F_{\Upsilon}^{(n)}(A) = \lim_{w_k \to -i\infty} \frac{\left\langle \prod_{k=0}^{n-1} \Upsilon(w_k) \Upsilon^{\dagger}(-w_k) \right\rangle_{\mathcal{R}_n}}{\langle \Upsilon(w_0) \Upsilon^{\dagger}(-w_0) \rangle_{\mathcal{R}_1}^n}$$

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Difference: the inserted operator is chiral!

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- Difference: the inserted operator is chiral!
- *R_n* is the replica manifold, where each Riemann sheet is now given by an infinite strip

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Entanglement entropies and boundary CFT (II)

 F can be simplified by means of the conformal transformations (R_n → D_n → D)

$$w \to \frac{\sin \frac{\pi(w-l)}{2L}}{\sin \frac{\pi(w+l)}{2L}} \to z = \left[\frac{\sin \frac{\pi(w-l)}{2L}}{\sin \frac{\pi(w+l)}{2L}}\right]^{\frac{1}{n}}$$

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► The operators are now placed on the boundary of the disk, at z[±]_{n,k}

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Entanglement entropies and boundary CFT (III)



Correlators on the disk become then correlators on C:

$$F_{\Upsilon}^{(n)}(A) = \frac{e^{i2\pi(n-1)h}}{n^{2nh}} \frac{\left\langle \prod_{k=0}^{n-1} \Upsilon(z_{n,k}^{-}) \Upsilon^{\dagger}(z_{n,k}^{+}) \right\rangle}{\left\langle \Upsilon_{0}(z_{1,0}^{-}) \Upsilon_{0}^{\dagger}(z_{1,0}^{+}) \right\rangle^{n}}$$

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In principle, we have a recipe to compute corrections!

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Spin-1/2 Ising chain:

$$H = -\frac{1}{2} \left(\sum_{j=1}^{L-1} \sigma_j^x \sigma_{j+1}^x + h \sum_{j=1}^{L} \sigma_j^z \right)$$

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Spin-1/2 Ising chain:

$$H = -\frac{1}{2} \left(\sum_{j=1}^{L-1} \sigma_j^x \sigma_{j+1}^x + h \sum_{j=1}^{L} \sigma_j^z \right)$$

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Critical for h = 1

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Spin-1/2 Ising chain:

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- Effective low-energy description: c = 1/2 minimal CFT
- Primary fields: $\mathbb{I}(h=0)$, $\sigma(h=\frac{1}{16})$, $\chi(h=\frac{1}{2})$

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Conformal boundary conditions: + (0), - (¹/₂) and F (¹/₁₆), corresponding to fixing σ[×] to +1/2, -1/2 or letting it free at the boundary [Cardy '86, Zhou *et al.* '06]

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- Fusion coefficients:

$$\mathcal{N}^{h} = \begin{array}{ccc} 0 \\ 1/16 \\ 1/2 \end{array} \begin{pmatrix} \delta^{h}_{0} & \delta^{h}_{1/16} & \delta^{h}_{1/2} \\ 0 & \delta^{h}_{0} + \delta^{h}_{1/2} & \delta^{h}_{1/16} \\ 0 & 0 & \delta^{h}_{0} \end{pmatrix} \\ 0 & 1/16 & 1/2 \end{array}$$

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- ► Example: Z_{1/16,1/16} = Z₀₀ = χ₀ + χ_{1/2} → the GS should get no corrections, while the first excited should originate from the primary χ
- ► All the needed correlation functions for the c = 1/2 minimal CFT are known [Ardonne and Sierra '10, Berganza et al. '12, Essler et al. '13] (see also F. Essler's talk)

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Critical Ising chain: numerics (I)

- For the FF case: free-fermions techniques [Vidal et al. '03, Peschel '03]
- In all other cases: DMRG simulations (thanks to Fabio Ortolani for providing the code)

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$$\lim_{n \to 1} \frac{1}{1-n} \ln F_{\chi}^{(n)}(x)$$

$$\ln |2 \sin(\pi x)| + \psi \left(\frac{1}{2 \sin(\pi x)}\right)$$

$$\sin(\pi x)$$

 $(S_{2,+F}-S_{2,FF})(l,L)/\ln 2$ $0.4 - (S_{\perp} - S_{\perp \perp})(l,L)/\ln 2$ I = 100I = 120I - 140L = 160(a) 0.3 L=180-0.2 CFT L = 600.2 -0.3 I = 80L = 100(a) L = 120-0.4 I = 1400.1 L = 160L = 180CFT х < □ > < □ > < □ > < □ > < □ > < E

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Critical Ising chain: numerics (II) Observations:

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Critical Ising chain: numerics (II)

Observations:

 Convergence to the CFT prediction increasing L, as confirmed by a FSS analysis: strong finite-size effect (see also [Igloi and Juhász '08])



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 Only thing we cannot predict: constant (and completely known) boundary entropies (BE's) [Affleck and Ludwig '91]: we add them to the CFT predictions

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- Only thing we cannot predict: constant (and completely known) boundary entropies (BE's) [Affleck and Ludwig '91]: we add them to the CFT predictions
- In any considered case, excellent agreement between CFT and numerics

 BC cannot be implemented "exactly" with DMRG, and one has to consider the modified Hamiltonian [Bilstein '00]:

$$\begin{aligned} H &= -\sum_{j=1}^{L-1} \left(\sigma_j^x \sigma_{j+1}^x + \sigma_j^y \sigma_{j+1}^y \right) + \\ &- \frac{1}{2} (\alpha_- \sigma_1^- + \alpha_+ \sigma_1^+ + \alpha_z \sigma_1^z + \beta_- \sigma_L^- + \beta_+ \sigma_L^+ + \beta_z \sigma_L^z) \end{aligned}$$

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Effective low-energy description: c = 1 CFT: compactified free massless boson Entanglement entropies in conformal systems with boundaries

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- Effective low-energy description: c = 1 CFT: compactified free massless boson
- From theory: possible conformal BC: Dirichlet (D) and Neumann (N)
- Lattice realization of conformal BC: D is realized by setting all boundary couplings to 0, N (on the left edge) by, e.g., α_−, α₊ ≠ 0

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The spin- $\frac{1}{2}$ XX chain (II)

Partition functions:

$$Z_{DD}(q) = K_0(q) + K_2(q) = q^{-\frac{1}{24}} \left[1 + 2q^{\frac{1}{2}} + q + O(q^2) \right]$$
$$Z_{NN}(q) = K_0(q) = q^{-\frac{1}{24}} \left[1 + q + O(q^2) \right]$$
$$Z_{ND}(q) = \chi_{\frac{1}{16}}(q) \left[\chi_0(q) + \chi_{\frac{1}{2}}(q) \right]$$

 K_j , $j = 0, 1, 2, 3 \mod 4$: characters of the c = 1 CFT(with unit compactification radius) [Saleur '98]

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- The operator content of the theories is, even in this case, under control
- Even in this case, the needed correlators are known

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 DD case: free-fermions techniques; remaining cases: DMRG Entanglement entropies in conformal systems with boundaries

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Observations:

Different lattice/finite-size effect with respect to the c = 1/2 case: strong parity oscillations, typical of c = 1 systems [Laflorencie et al. '06, Calabrese et al. '10, Dalmonte, Ercolessi & LT '11, '12] Entanglement entropies in conformal systems with boundaries

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- Excellent agreement in any considered case

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- Starting from the results for excess entropies of excited states, we derived a CFT master formula for the excess entropies in the open cases
- Starting from it, computations are straightforward once one knows how to compute *n*-point primary correlations

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- ► Message: boundaries reduce the operator content of CFT → open models are advantageous playgrounds for the testing of CFT predictions
- Example of possible open issue: deriving the form of the corrections for a generic descendant state in CFT (work in progress!) and of the oscillating corrections

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Thank you for the attention!

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